

# Matrix Formulation of the Subsidiary Conditions of the Second Kind for the Pauli-Fierz Wave Equation

C. G. Koutroulos

Department of Theoretical Physics, University of Thessaloniki, Thessaloniki, Greece

Z. Naturforsch. **41a**, 1029–1030 (1986); received March 18, 1986

It is demonstrated, using the Pauli-Fierz wave equation as an example, that the subsidiary conditions of the second kind for a first order wave equation can be expressed in matrix language.

In a series of letters [1, 2, 3, 4] we saw how the subsidiary conditions of the second kind (which are important in studying the propagation of a wave equation in an external field) can be verified by employing spinor calculus. This method has the disadvantage that it is rather complicated. Thus in this paper we show how the same problem can be formulated in matrix language. We use as a prototype the Pauli-Fierz wave equation [5] and show that its subsidiary conditions of the second kind can be expressed in matrix form. In a subsequent paper we shall derive general conditions which have to be satisfied such that the subsidiary conditions of the second kind for any first order wave equation be derivable.

For the moment let us consider the Pauli-Fierz wave equation given in [1, 2]. As it was shown there, this wave equation accepts the following subsidiary conditions of the second kind:

$$6\kappa^2 d^{\dot{1}} + (f_{10} + f_{13} + if_{02} + if_{32}) a_{12}^{\dot{1}} + 2(f_{03} + if_{12}) a_{12}^{\dot{1}} + (f_{01} + if_{02} + f_{13} + if_{23}) \alpha_{22}^{\dot{1}} + 2(f_{03} - if_{12}) d^{\dot{1}} + 2(f_{01} + if_{20} - f_{13} - if_{23}) d^{\dot{2}} = 0, \quad (1)$$

$$6\kappa^2 d^{\dot{2}} + (-f_{01} - if_{20} + f_{13} + if_{32}) a_{11}^{\dot{2}} + 2(-f_{30} + if_{12}) a_{12}^{\dot{2}} + (-f_{10} + f_{13} - if_{20} + if_{23}) a_{22}^{\dot{2}} + 2(f_{30} - if_{21}) d^{\dot{2}} + 2(f_{01} - f_{31} + if_{20} - if_{23}) d^{\dot{1}} = 0, \quad (2)$$

$$6\kappa^2 c_1 + (-f_{01} + f_{31} - if_{20} + if_{23}) b_{11}^{\dot{1}} + 2(if_{21} - f_{30}) b_{11}^{\dot{2}} + (-f_{10} - if_{20} + f_{31} + if_{32}) b_{11}^{\dot{2}} + 2(f_{30} - if_{12}) c_1 + 2(f_{10} + if_{20} - f_{13} - if_{23}) c_2 = 0, \quad (3)$$

Reprint requests to Dr. C. G. Koutroulos, Department of Theoretical Physics, University of Thessaloniki, Thessaloniki 54006, Griechenland.

$$6\kappa^2 c_2 + (-f_{01} + f_{31} - if_{20} + if_{23}) b_{21}^{\dot{1}} + 2(-f_{30} + if_{21}) b_{21}^{\dot{2}} + (-f_{10} + f_{31} - if_{20} + if_{32}) b_{21}^{\dot{2}} + 2(f_{10} - f_{31} + if_{02} - if_{23}) c_1 + 2(-if_{21} + f_{03}) c_2 = 0. \quad (4)$$

Let us for instance consider the derivation of the subsidiary condition (1). In deriving this subsidiary condition the following steps were involved (see [2]). First construct

$$\{\pi_1^{\dot{1}}(eq_9) + \pi_2^{\dot{1}}(eq_{10}) + \pi_1^{\dot{2}}(eq_{12}) + \pi_2^{\dot{2}}(eq_{13})\} - \{3\pi_1^{\dot{1}}(eq_{15}) - 3\pi_2^{\dot{2}}(eq_{16})\} = 0 \quad (5)$$

and substitute into it  $6\kappa^6(eq_7)$  having replaced  $\pi_{a\dot{a}}$  to produce terms involving  $[\pi_k, \pi_l] = i e F_{kl} = f_{kl}$ .

The above steps can be put into matrix language as follows:

$$\mathbf{S}^{(I)} \{\mathbf{A}_0^{(I)} \pi_0 + \mathbf{A}_1^{(I)} \pi_1 + \mathbf{A}_2^{(I)} p_2 + \mathbf{A}_3^{(I)} \pi_3 + \mathbf{A}_4^{(I)} \kappa\} \cdot \{-\mathbb{L}_0 \pi_0 + \mathbb{L}_1 \pi_1 + \mathbb{L}_2 \pi_2 + \mathbb{L}_3 \pi_3 + \mathbb{1} \kappa\} \psi = 0, \quad (6)$$

where  $\mathbf{S}^{(I)}$  is an  $1 \times 16$  matrix and  $\mathbf{A}_0^{(I)}, \mathbf{A}_1^{(I)}, \mathbf{A}_2^{(I)}, \mathbf{A}_3^{(I)}, \mathbf{A}_4^{(I)}$  are  $16 \times 16$  matrices. The index (I) makes references to the first subsidiary condition of the second kind. The matrices  $\mathbf{S}^{(I)}, \mathbf{A}_j^{(I)}, j = 0, 1, 2, 3, 4$ , with respect to the spinor basis

$$\{a_{11}^{\dot{1}}, a_{12}^{\dot{1}}, a_{22}^{\dot{1}}, a_{11}^{\dot{2}}, a_{12}^{\dot{2}}, a_{22}^{\dot{2}}, d^{\dot{1}}, d^{\dot{2}}, b_{11}^{\dot{1}}, b_{12}^{\dot{1}}, b_{22}^{\dot{1}}, b_{11}^{\dot{2}}, b_{12}^{\dot{2}}, b_{22}^{\dot{2}}, c_1, c_2\} \quad (7)$$

are

$$\mathbf{S}^{(I)} = [0, 0, 0, 0, 0, 0, 6, 0, 1, 1, 0, 1, 1, 0, -3, -3], \quad (8)$$

$$\mathbf{A}_0^{(I)} = \text{diag}[0, 0, 0, 0, 0, 0, 0, 0, -1, 0, 1, 0, 0, -1, 0], \quad (9)$$

$$\mathbf{A}_1^{(I)} = \text{diag}[0, 0, 0, 0, 0, 0, 0, 0, 1, 0, 0, 0, -1, 0, 0, -1], \quad (10)$$

$$\mathbf{A}_2^{(I)} = \text{diag}[0, 0, 0, 0, 0, 0, 0, 0, -i, 0, 0, 0, -i, 0, 0, -i], \quad (11)$$

$$\mathbf{A}_3^{(I)} = \text{diag}[0, 0, 0, 0, 0, 0, 0, 0, -1, 0, -1, 0, 0, -1, 0], \quad (12)$$

$$\mathbf{A}_4^{(I)} = \text{diag}[0, 0, 0, 0, 0, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0]. \quad (13)$$

0340-4811 / 86 / 0800-1029 \$ 01.30/0. – Please order a reprint rather than making your own copy.



Dieses Werk wurde im Jahr 2013 vom Verlag Zeitschrift für Naturforschung in Zusammenarbeit mit der Max-Planck-Gesellschaft zur Förderung der Wissenschaften e.V. digitalisiert und unter folgender Lizenz veröffentlicht: Creative Commons Namensnennung-Keine Bearbeitung 3.0 Deutschland Lizenz.

Zum 01.01.2015 ist eine Anpassung der Lizenzbedingungen (Entfall der Creative Commons Lizenzbedingung „Keine Bearbeitung“) beabsichtigt, um eine Nachnutzung auch im Rahmen zukünftiger wissenschaftlicher Nutzungsformen zu ermöglichen.

This work has been digitalized and published in 2013 by Verlag Zeitschrift für Naturforschung in cooperation with the Max Planck Society for the Advancement of Science under a Creative Commons Attribution-NoDerivs 3.0 Germany License.

On 01.01.2015 it is planned to change the License Conditions (the removal of the Creative Commons License condition “no derivative works”). This is to allow reuse in the area of future scientific usage.

The matrices for the other three subsidiary conditions are, respectively,

$$\mathbf{S}^{(II)} = [0, 0, 0, 0, 0, 0, 6, 0, 1, 1, 0, 1, 1, -3, -3], \quad (14)$$

$$\mathbf{A}_0^{(II)} = \text{diag}[0, 0, 0, 0, 0, 0, 0, 0, 0, -1, 0, 1, 0, 0, -1], \quad (15)$$

$$\mathbf{A}_1^{(II)} = \text{diag}[0, 0, 0, 0, 0, 0, 0, 0, 0, 1, 0, 0, 0, -1, -1, 0], \quad (16)$$

$$\mathbf{A}_2^{(II)} = \text{diag}[0, 0, 0, 0, 0, 0, 0, 0, 0, -i, 0, 0, 0, -i, i, 0], \quad (17)$$

$$\mathbf{A}_3^{(II)} = \text{diag}[0, 0, 0, 0, 0, 0, 0, 0, 0, -1, 0, -1, 0, 0, 1], \quad (18)$$

$$\mathbf{A}_4^{(II)} = \text{diag}[0, 0, 0, 0, 0, 0, 0, 1, 0, 0, 0, 0, 0, 0, 0], \quad (19)$$

$$\mathbf{S}^{(III)} = [1, 1, 0, 1, 1, 0, -3, -3, 0, 0, 0, 0, 0, 0, 6, 0], \quad (20)$$

$$\mathbf{A}_0^{(III)} = \text{diag}[0, 1, 0, -1, 0, 0, -1, 0, 0, 0, 0, 0, 0, 0, 0, 0], \quad (21)$$

$$\mathbf{A}_1^{(III)} = \text{diag}[1, 0, 0, 0, -1, 0, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0], \quad (22)$$

$$\mathbf{A}_2^{(III)} = \text{diag}[-i, 0, 0, 0, -i, 0, 0, i, 0, 0, 0, 0, 0, 0, 0, 0], \quad (23)$$

$$\mathbf{A}_3^{(III)} = \text{diag}[0, -1, 0, -1, 0, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0], \quad (24)$$

$$\mathbf{A}_4^{(III)} = \text{diag}[0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 1, 0], \quad (25)$$

$$\mathbf{S}^{(IV)} = [0, 1, 1, 0, 1, 1, -3, -3, 0, 0, 0, 0, 0, 0, 0, 6], \quad (26)$$

$$\mathbf{A}_0^{(IV)} = \text{diag}[0, 0, 1, 0, -1, 0, 0, -1, 0, 0, 0, 0, 0, 0, 0, 0], \quad (27)$$

$$\mathbf{A}_1^{(IV)} = \text{diag}[1, 0, 0, 0, -1, 0, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0], \quad (28)$$

$$\mathbf{A}_2^{(IV)} = \text{diag}[-i, 0, 0, 0, -i, 0, 0, i, 0, 0, 0, 0, 0, 0, 0, 0], \quad (29)$$

$$\mathbf{A}_3^{(IV)} = \text{diag}[0, -1, 0, -1, 0, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0], \quad (30)$$

$$\mathbf{A}_4^{(IV)} = \text{diag}[0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 1], \quad (31)$$

If  $\mathbf{S}$  is considered as a  $4 \times 20$  matrix and  $\mathbf{A}_0, \mathbf{A}_1, \mathbf{A}_2, \mathbf{A}_3, \mathbf{A}_4$  are considered as  $20 \times 16$  matrices then all four subsidiary conditions can be accommodated into one matrix formula, namely

$$\mathbf{S} \{ \mathbf{A}_0 \pi_0 + \mathbf{A}_1 \pi_1 + \mathbf{A}_2 \pi_2 + \mathbf{A}_3 \pi_3 + \mathbf{A}_4 \pi_4 \} \\ \cdot \{ -\mathbb{L}_0 \pi_0 + \mathbb{L}_1 \pi_1 + \mathbb{L}_2 \pi_2 + \mathbb{L}_3 \pi_3 + \mathbb{L}_4 \pi_4 \} \psi = 0.$$

Thus we have demonstrated with this example that the subsidiary conditions of the second kind of a first order wave equation can be derived using matrix calculus. The question which arises is what conditions must be satisfied such that the subsidiary conditions of the second kind for any first order wave equation be derivable. We shall deal with this question in a subsequent paper.

- [1] C. G. Koutroulos, Lett. Nuovo Cim. **39**, 96 (1984).
- [2] C. G. Koutroulos, Lett. Nuovo Cim. **38**, 217 (1983).
- [3] C. G. Koutroulos, Lett. Nuovo Cim. **40**, 339 (1984).

- [4] C. G. Koutroulos, Lett. Nuovo Cim. **41**, 112 (1984).
- [5] M. Fierz and W. Pauli, Proc. Roy. Soc. London, Ser. A **173**, 211 (1939).